

THE ELECTROSTATIC LIMIT FOR THE 3D ZAKHAROV SYSTEM

PAOLO ANTONELLI AND LUIGI FORCELLA

ABSTRACT. We consider the vectorial Zakharov system describing Langmuir waves in a weakly magnetized plasma. In its original derivation [Z] the evolution for the electric field envelope is governed by a Schrödinger type equation with a singular parameter which is usually large in physical applications. Motivated by this, we study the rigorous limit as this parameter goes to infinity. By using some Strichartz type estimates to control separately the fast and slow dynamics in the problem, we show that the evolution of the electric field envelope is asymptotically constrained onto the space of irrotational vector fields.

1. INTRODUCTION.

In this paper we consider the vectorial Zakharov system [Z] describing Langmuir waves in a weakly magnetized plasma. After a suitable rescaling of the variables it reads [SS]

$$(1.1) \quad \begin{cases} i\partial_t u - \alpha \nabla \wedge \nabla \wedge u + \nabla(\operatorname{div} u) = nu \\ \frac{1}{c_s^2} \partial_{tt} n - \Delta n = \Delta |u|^2 \end{cases} ,$$

subject to initial conditions

$$u(0) = u_0, \quad n(0) = n_0, \quad \partial_t n(0) = n_1.$$

Here $u : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{C}^3$ describes the slowly varying envelope of the highly oscillating electric field, whereas $n : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}$ is the ion density fluctuation. The rescaled constants in (1.1) are $\alpha = \frac{c^2}{3v_e^2}$, $v_e = \sqrt{\frac{T_e}{m_e}}$ being the electron thermal velocity, while c_s is proportional to the ion acoustic speed. In many physical situations the parameter α is relatively large, see for example table 1, p. 47 in [TtH]. In this regime the electric field is almost irrotational and in the electrostatic limit $\alpha \rightarrow \infty$ the dynamics is asymptotically described by

$$(1.2) \quad \begin{cases} i\partial_t u + \Delta u = \mathbf{Q}(nu) \\ \frac{1}{c_s^2} \partial_{tt} n - \Delta n = \Delta |u|^2 \end{cases} ,$$

where $\mathbf{Q} = -(-\Delta)^{-1} \nabla \operatorname{div}$ is the Helmholtz projection operator onto irrotational vector fields. By further simplifying (1.1) it is possible to consider the so called scalar Zakharov system

$$(1.3) \quad \begin{cases} i\partial_t u + \Delta u = nu \\ \frac{1}{c_s^2} \partial_{tt} n - \Delta n = \Delta |u|^2 \end{cases} ,$$

which retains the main features of (1.2). In the subsonic limit $c_s \rightarrow \infty$ we find the cubic focusing nonlinear Schrödinger equation

$$i\partial_t u + \Delta u + |u|^2 u = 0.$$

The Cauchy problem for the Zakharov system has been extensively studied in the mathematical literature. For the local and global well-posedness, see [SS2, OT1, OT2, KPV, BC] and the recent results concerning low regularity solutions [GTV, BH]. In [M] formation of blow-up solutions is studied by means of virial identities, see also [GM] where self-similar solutions are constructed in two space dimensions. The subsonic limit $c_s \rightarrow \infty$ for (1.3) is investigated in [SW]. Furthermore, some related singular limits are also studied in [MN], considering the Klein-Gordon-Zakharov system.

The aim of this paper is to rigorously study the electrostatic limit for the vectorial Zakharov equation, namely we show that mild solutions to (1.1) converge towards solutions to (1.2) as $\alpha \rightarrow \infty$.

As we will see below, we will investigate this limit by exploiting two auxiliary systems associated to (1.1), (1.2), namely systems (3.1) and (4.1) below. This approach, already introduced in [OT1, OT2] to study local and global well-posedness for the Zakharov system (1.3), overcomes the problem generated by the loss of derivatives on the term $|u|^2$ in the wave equation, but in our context it introduces a new difficulty. Indeed the initial data for the auxiliary system (3.1) are not uniformly bounded for $\alpha \geq 1$, see for example (4.2) and the discussion at the beginning of Section 4 below.

For this reason we will need to consider a family of well-prepared initial data; more precisely we will take a set u_0^α of initial states for the Schrödinger part in (1.1) which converges to an irrotational initial datum for (1.2).

It is reasonable to think that by using similar arguments to [BH] then it is possible to study the electrostatic limit for rougher solutions to (1.1), in order to skip the assumption on the well-prepared data. For example, in [G] an analogous limit for a nonlinear Schrödinger type equation is performed, at the regularity level of L^2 solutions; in that paper the convergence holds true for any initial datum. Let us notice however that the equation studied in [G] is of nonlinear Schrödinger type, and so it does not experience the phenomenon of derivative loss as for (1.1). The problem of studying rough solutions to (1.1) and their electrostatic limit, although being an interesting and physically relevant result, goes beyond the scope of this paper and it could be subject of future research.

We consider initial data $(u_0^\alpha, n_0^\alpha, n_1^\alpha) \in H^2(\mathbb{R}^3) \times H^1(\mathbb{R}^3) \times L^2(\mathbb{R}^3) =: \mathcal{H}_2$ for (1.1), converging in the same space to a set of initial data $(u_0^\infty, n_0^\infty, n_1^\infty) \in \mathcal{H}_2$, with u_0^∞ an irrotational vector field, and we show the convergence in the space

$$\begin{aligned} \mathcal{X}_T := \{ (u, n) : u \in L^q(0, T; W^{2,r}(\mathbb{R}^3)), \forall (q, r) \text{ admissible pair,} \\ n \in L^\infty(0, T; H^1(\mathbb{R}^3)) \cap W^{1,\infty}(0, T; L^2(\mathbb{R}^3)) \}. \end{aligned}$$

For a more detailed discussion about notations and the spaces considered in this paper we refer the reader to Section 2.

Before stating our main result we first recall the local well-posedness result in \mathcal{H}_2 for system (1.2).

Theorem 1.1 ([OT1]). *Let $(u_0, n_0, n_1) \in \mathcal{H}_2$, then there exist a maximal time $0 < T_{max} \leq \infty$ and a unique solution (u, n) to (1.2) such that $u \in \mathcal{C}([0, T_{max}); H^2) \cap$*

$\mathcal{C}^1([0, T_{max}); L^2)$, $n \in \mathcal{C}([0, T_{max}); H^1) \cap \mathcal{C}^1([0, T_{max}); L^2)$. Furthermore the solution depends continuously on the initial data and the standard blow-up alternative holds true: either $T_{max} = \infty$ and the solution is global or $T_{max} < \infty$ and we have

$$\lim_{t \rightarrow T_{max}} \|(u, n, \partial_t n)(t)\|_{\mathcal{H}_2} = \infty.$$

Analogously we are going to prove the same local well-posedness result for system (1.1). Moreover, despite of the fact that the initial datum for (3.1) is not uniformly bounded for $\alpha \geq 1$ (see the discussion at the beginning of Section 3), we can anyway infer some a priori bounds in α for the solution (u^α, n^α) to (1.1).

Theorem 1.2. *Let $(u_0^\alpha, n_0^\alpha, n_1^\alpha) \in \mathcal{H}_2$, then there exist a maximal time $T_{max}^\alpha > 0$ and a unique solution (u^α, n^α) to (1.1) such that*

- $u^\alpha \in \mathcal{C}([0, T_{max}^\alpha); H^2) \cap \mathcal{C}^1([0, T_{max}^\alpha); L^2)$,
- $n^\alpha \in \mathcal{C}([0, T_{max}^\alpha); H^1) \cap \mathcal{C}^1([0, T_{max}^\alpha); L^2)$.

Furthermore the existence times T_{max}^α are uniformly bounded from below, $0 < T^* \leq T_{max}^\alpha$ for any $\alpha \geq 1$, and we have

$$\|(u^\alpha, n^\alpha, \partial_t n^\alpha)\|_{L^\infty(0, T; \mathcal{H}_2)} + \|\partial_t u^\alpha\|_{L^2(0, T; L^6)} \leq C(T, M),$$

for any $0 < T < T_{max}^\alpha$, where the constant above does not depend on $\alpha \geq 1$.

Our main result in this paper is the following one.

Theorem 1.3. *Let $(u_0^\alpha, n_0^\alpha, n_1^\alpha) \in \mathcal{H}_2$ and let (u^α, n^α) be the maximal solution to (1.1) defined on the time interval $[0, T_{max}^\alpha)$. Let us assume that*

$$\lim_{\alpha \rightarrow \infty} \|(u_0^\alpha, n_0^\alpha, n_1^\alpha) - (u_0^\infty, n_0^\infty, n_1^\infty)\|_{\mathcal{H}_2} = 0,$$

for some $(u_0^\infty, n_0^\infty, n_1^\infty) \in \mathcal{H}_2$ such that $u_0^\infty = \mathbf{Q}u_0^\infty$, and let (u^∞, n^∞) be the maximal solutions to (1.2) in the interval $[0, T_{max}^\infty)$ with such initial data. Then

$$\liminf_{\alpha \rightarrow \infty} T_{max}^\alpha \geq T_{max}^\infty$$

and we have the following convergence

$$\lim_{\alpha \rightarrow \infty} \|(u^\alpha, n^\alpha) - (u^\infty, n^\infty)\|_{\mathcal{X}_T} = 0,$$

for any $0 < T < T_{max}^\infty$.

The paper is structured as follows. In Section 2 we fix some notations and give some preliminary results which will be used in the analysis of the problem below. In Section 3 we show the local well-posedness of system (1.1) in the space \mathcal{H}_2 . Finally in Section 4 we investigate the electrostatic limit and prove the main theorem.

Acknowledgement. This paper and its project originated after many useful discussions with Prof. Pierangelo Marcati, during second author's M. Sc. thesis work. We would like to thank P. Marcati for valuable suggestions.

2. PRELIMINARY RESULTS AND TOOLS.

In this section we introduce notations and some preliminary results which will be useful in the analysis below. The Fourier transform of a function f is defined by

$$\mathcal{F}(f)(\xi) = \hat{f}(\xi) = \int_{\mathbb{R}^3} e^{-2\pi i x \cdot \xi} f(x) dx,$$

with its inverse

$$f(x) = \int_{\mathbb{R}^3} e^{2\pi i x \cdot \xi} \hat{f}(\xi) d\xi.$$

Given an interval $I \subset \mathbb{R}$, we denote by $L^q(I; L^r)$ the Bochner space equipped with the norm defined by

$$\|f\|_{L^q(I; L^r)} = \left(\int_I \|f(s)\|_{L^r(\mathbb{R}^3)}^q ds \right)^{1/q},$$

where $f = f(s, x)$. When no confusion is possible, we write $L_t^q L_x^r = L^q(I; L^r(\mathbb{R}^3))$. Given two Banach spaces X, Y , we denote $\|f\|_{X \cap Y} := \max\{\|f\|_X, \|f\|_Y\}$ for $f \in X \cap Y$. With $W^{k,p}$ we denote the standard Sobolev spaces and for $p = 2$ we write $H^k = W^{k,2}$. $A \lesssim B$ means that there exists a universal constant C such that $A \leq CB$ and in general in a chain of inequalities the constant may change from one line to the other.

As already said in the Introduction, given a vector field F , we denote by $\mathbf{Q}F = -(-\Delta)^{-1} \nabla \operatorname{div} F$ its projection into irrotational fields, moreover $\mathbf{P} = 1 - \mathbf{Q}$ is its orthogonal projection operator onto solenoidal fields.

The space of initial data is denoted by $\mathcal{H}_2 := H^2(\mathbb{R}^3) \times H^1(\mathbb{R}^3) \times L^2(\mathbb{R}^3)$. A pair of Lebesgue exponents is called *Schrödinger admissible* (or simply admissible) if $2 \leq q \leq \infty$, $2 \leq r \leq 6$ and they are related through

$$\frac{1}{q} = \frac{3}{2} \left(\frac{1}{2} - \frac{1}{r} \right).$$

Given a time interval $I \subset \mathbb{R}$ we denote the Strichartz space $S^0(I)$ to be the closure of the Schwartz space with the norm

$$\|u\|_{S^0(I)} := \sup_{(q,r)} \|u\|_{L^q(I; L^r(\mathbb{R}^3))},$$

where the sup is taken over all admissible pairs; furthermore we write

$$S^2(I) = \{u \in S^0(I) : \nabla^2 u \in S^0(I)\}.$$

We define moreover the space

$$\mathcal{W}^1(I) = \{n : n \in L^\infty(I; H^1) \cap W^{1,\infty}(I; L^2)\}$$

endowed with the norm

$$\|n\|_{\mathcal{W}^1(I)} = \|n\|_{L^\infty(I; H^1)} + \|\partial_t n(t)\|_{L^\infty(I; L^2)}.$$

The space of solutions we consider in this paper is given by

$$\mathcal{X}_T = \{(u, n) : u \in S^2([0, T]), n \in \mathcal{W}^1([0, T])\}.$$

We will also use the following notation:

$$\begin{aligned} \mathcal{C}([0, T]; \mathcal{H}_2) = \{ & (u, n) : u \in \mathcal{C}([0, T]; H^2) \cap \mathcal{C}^1([0, T]; L^2), \\ & n \in \mathcal{C}([0, T]; H^1) \cap \mathcal{C}^1([0, T]; L^2) \}. \end{aligned}$$

Here in this paper we only consider positive times, however the same results are valid also for negative times.

We now introduce some basic preliminary results which will be useful later in the analysis.

First of all we consider the linear propagator related to (1.1), namely

$$(2.1) \quad i\partial_t u = \alpha \nabla \wedge \nabla \wedge u - \nabla \operatorname{div} u.$$

Lemma 2.1. *Let u solve (2.1) with initial datum $u(0) = u_0$, then*

$$(2.2) \quad u(t) = U_Z(t)u_0 = [U(\alpha t)\mathbf{P} + U(t)\mathbf{Q}]u_0,$$

where $U(t) = e^{it\Delta}$ is the Schrödinger evolution operator.

Proof. By taking the Fourier transform (2.1) we have

$$\begin{aligned} i\partial_t \hat{u} &= -\alpha\xi \wedge \xi \wedge \hat{u} + \xi(\xi \cdot \hat{u}) \\ &= |\xi|^2 \left(\alpha\hat{\mathbf{P}}(\xi) + \hat{\mathbf{Q}}(\xi) \right) \hat{u}(\xi), \end{aligned}$$

where $\hat{\mathbf{P}}(\xi), \hat{\mathbf{Q}}(\xi)$ are two (3×3) -matrices defined by $\hat{\mathbf{Q}}(\xi) = \frac{\xi \otimes \xi}{|\xi|^2}$, $\hat{\mathbf{P}}(\xi) = \mathbf{1} - \hat{\mathbf{Q}}(\xi)$ where $\mathbf{1}$ is the identity matrix. Hence we may write

$$\hat{u}(t) = e^{-i\alpha t|\xi|^2 \hat{\mathbf{P}}(\xi) - it|\xi|^2 \hat{\mathbf{Q}}(\xi)} \hat{u}_0(\xi).$$

It is straightforward to see that $\hat{\mathbf{Q}}(\xi)$ is a projection matrix, $0 \leq \hat{\mathbf{Q}}(\xi) \leq 1$, $\hat{\mathbf{Q}}(\xi) = \hat{\mathbf{Q}}^2(\xi)$, hence $\hat{\mathbf{P}}(\xi)$ is its orthogonal projection. Consequently we have

$$\begin{aligned} \hat{u}(t) &= e^{-i\alpha t|\xi|^2 \hat{\mathbf{P}}(\xi)} e^{-it|\xi|^2 \hat{\mathbf{Q}}(\xi)} \hat{u}_0(\xi) \\ &= \left(e^{-i\alpha t|\xi|^2} \hat{\mathbf{P}}(\xi) + \hat{\mathbf{Q}}(\xi) \right) \left(e^{-it|\xi|^2} \hat{\mathbf{Q}}(\xi) + \hat{\mathbf{P}}(\xi) \right) \hat{u}_0(\xi) \\ &= \left(e^{-i\alpha t|\xi|^2} \hat{\mathbf{P}}(\xi) + e^{-it|\xi|^2} \hat{\mathbf{Q}}(\xi) \right) \hat{u}_0(\xi). \end{aligned}$$

By taking the inverse Fourier transform we find (2.2). \square

By the dispersive estimates for the standard Schrödinger evolution operator we have

$$(2.3) \quad \begin{aligned} \|U(t)\mathbf{Q}f\|_{L^p} &\lesssim |t|^{-3(\frac{1}{2}-\frac{1}{p})} \|\mathbf{Q}f\|_{L^{p'}}, \\ \|U(\alpha t)\mathbf{P}f\|_{L^p} &\lesssim |\alpha t|^{-3(\frac{1}{2}-\frac{1}{p})} \|\mathbf{P}f\|_{L^{p'}}, \end{aligned}$$

for any $2 \leq p \leq \infty$, $t \neq 0$. These two estimates together give

$$\|U_Z(t)f\|_{L_x^p} \lesssim |t|^{-3(\frac{1}{2}-\frac{1}{p})} \|f\|_{L_x^{p'}},$$

for $2 \leq p < \infty$. Let us notice that the dispersive estimate for $p = \infty$ does not hold for $U_Z(t)$ anymore because the projection operators \mathbf{Q}, \mathbf{P} are not bounded from L^1 into itself. Nevertheless by using the dispersive estimates in (2.3) and the result in [KT] we infer the whole set of Strichartz estimates for the irrotational and solenoidal part, separately. By summing them up we thus find the Strichartz estimates for the propagator in (2.2).

Lemma 2.2. *Let $(q, r), (\gamma, \rho)$ be two arbitrary admissible pairs and let $\alpha \geq 1$, then we have*

$$(2.4) \quad \begin{aligned} \|U(\alpha t)\mathbf{P}f\|_{L_t^q(I; L_x^r)} &\leq C\alpha^{-\frac{2}{q}} \|f\|_{L_x^2}, \\ \left\| \int_0^t U(\alpha(t-s))\mathbf{P}F(s) ds \right\|_{L_t^q(I; L_x^r)} &\leq C\alpha^{-(\frac{1}{q}+\frac{1}{\gamma})} \|F\|_{L_t^{\gamma'}(I; L_x^{\rho'})}. \end{aligned}$$

and

$$\begin{aligned} \|U(t)\mathbf{Q}f\|_{L_t^q(I; L_x^r)} &\leq C \|f\|_{L_x^2}, \\ \left\| \int_0^t U(t-s)\mathbf{Q}F(s) ds \right\|_{L_t^q(I; L_x^r)} &\leq C \|F\|_{L_t^{\gamma'}(I; L_x^{\rho'})}, \end{aligned}$$

Consequently we also have

$$(2.5) \quad \|U_Z(t)g\|_{L_t^q(I;L_x^r)} \leq C\|f\|_{L_x^2},$$

$$(2.6) \quad \left\| \int_0^t U_Z(t-s)F(s) ds \right\|_{L_t^q(I;L_x^r)} \leq C\|F\|_{L_t^{\gamma'}(I;L_x^{\rho'})}.$$

Remark 2.3. • From the estimates in the Lemma above it is already straightforward that, at least in the linear evolution, we can separate the fast and slow dynamics and that the fast one is asymptotically vanishing. This is somehow similar to what happens with rapidly varying dispersion management, see for example [ASS].

- Let us notice that the constants in (2.5) and (2.6) are uniformly bounded for $\alpha \geq 1$. This is straightforward but it is a necessary remark to infer that the existence time in the local well-posedness section is uniformly bounded from below for any $\alpha \geq 1$.

3. LOCAL EXISTENCE THEORY.

In this Section we study the local well-posedness of (1.1) in the space \mathcal{H}_2 . We are going to perform a fixed point argument in order to find a unique local solution in the time interval $[0, T]$, for some $0 < T < \infty$. By standard arguments it is then possible to extend the solution up to a maximal time T_{max} for which the blow-up alternative holds. However, due to the loss of derivatives on the term $|u|^2$, we cannot proceed in a straightforward way, thus we follow the approach in [OT1] where the authors use an auxiliary system to overcome this difficulty. More precisely, let us define $v := \partial_t u$, then by differentiating the Schrödinger equation in (1.1) with respect to time, we write the following system

$$(3.1) \quad \begin{cases} i\partial_t v - \alpha \nabla \wedge \nabla \wedge v + \nabla \operatorname{div} v = nv + \partial_t nu \\ \partial_{tt} n - \Delta n = \Delta |u|^2 \\ iv - \alpha \nabla \wedge \nabla \wedge u + \nabla \operatorname{div} u = nu \end{cases}.$$

Differently from [OT1], here we encounter a further difficulty. Indeed we have that the initial datum for v is given by

$$(3.2) \quad v(0) = -i\alpha \nabla \wedge \nabla \wedge u_0 + i\nabla \operatorname{div} u_0 - in_0 u_0,$$

which in general is not uniformly bounded in L^2 for $\alpha \geq 1$. Hence the standard fixed point argument applied to the integral formulation of (3.1) would give a local solution on a time interval $[0, T^\alpha]$, where T^α goes to zero as α goes to infinity. For this reason we introduce the alternative variable

$$\tilde{v}(t) := v(t) - U(\alpha t)\mathbf{P}(i\alpha \Delta u_0),$$

for which we prove that the existence time T^α is uniformly bounded from below for $\alpha \geq 1$. The main result of this Section concerns the local well-posedness for (3.1).

Proposition 3.1. *Let $(u_0, n_0, n_1) \in \mathcal{H}_2$ be such that*

$$M := \|(u_0, n_0, n_1)\|_{\mathcal{H}_2}.$$

Then, for any $\alpha \geq 1$ there exists $\tau = \tau(M)$ and a unique local solution $(u, n) \in \mathcal{C}([0, \tau]; \mathcal{H}_2)$ to (1.1) such that

$$\sup_{[0, \tau]} \|(u, n, \partial_t n)(t)\|_{\mathcal{H}_2} \leq 2M$$

and

$$\|v\|_{L_t^2 L_x^6} \leq CM,$$

where C does not depend on $\alpha \geq 1$.

By standard arguments we then extend the local solution in [Proposition 3.1](#) to a maximal existence interval where the standard blow-up alternative holds true.

Theorem 3.2. *Let $(u_0, n_0, n_1) \in \mathcal{H}_2$, then for any $\alpha \geq 1$ there exists a unique maximal solution $(u^\alpha, v^\alpha, n^\alpha)$ to (3.1) with initial data $(u_0, v(0), n_0, n_1)$, $v(0)$ given by (3.2), on the maximal existence interval $I_\alpha := [0, T_{max}^\alpha)$, for some $T_{max}^\alpha > 0$. The solution satisfies the following regularity properties:*

- $u^\alpha \in \mathcal{C}(I_\alpha; H^2)$, $u^\alpha \in S^2([0, T])$, $\forall 0 < T < T_{max}^\alpha$,
- $v^\alpha \in \mathcal{C}(I_\alpha; L^2)$, $v^\alpha \in S^0([0, T])$, $\forall 0 < T < T_{max}^\alpha$,
- $n^\alpha \in \mathcal{C}(I_\alpha; H^1) \cap \mathcal{C}^1(I_\alpha; L^2)$.

Moreover, the following blow-up alternative holds true: $T_{max}^\alpha < \infty$ if and only if

$$\lim_{t \rightarrow T^\alpha} \|(u^\alpha(t), n^\alpha(t))\|_{\mathcal{H}_2} = \infty.$$

Finally, the map $\mathcal{H}_2 \rightarrow \mathcal{C}([0, T_{max}^\alpha); \mathcal{H}_2)$ associating any initial datum to its solution is a continuous operator.

Remark 3.3. The blow-up alternative above also implies in particular that the family of maximal existence times T^α is strictly bounded away from zero, i.e. there exists a $T^* > 0$ such that $T^* \leq T^\alpha$ for any $\alpha \geq 1$.

[Theorem 1.2](#) yields in a straightforward way from [Theorem 3.2](#) above.

Proof of Theorem 1.2. Let $(u^\alpha, v^\alpha, n^\alpha)$ be the solution to (3.1) constructed in [Theorem 3.2](#), then to prove the [Theorem 1.2](#) we only need to show that we identify $\partial_t u^\alpha = v^\alpha$ in the distributional sense. Let us differentiate with respect to t the equation

$$(1 - \alpha \Delta \mathbf{P} - \Delta \mathbf{Q})u = iv - (n - 1) \left(u_0 + \int_0^t v(s) ds \right)$$

obtaining

$$(3.3) \quad (1 - \alpha \Delta \mathbf{P} - \Delta \mathbf{Q})\partial_t u = i\partial_t v - (n - 1)v - \partial_t n \left(u_0 + \int_0^t v(s) ds \right),$$

this equation holding in H^{-2} , while the first equation of (3.1) gives us

$$(1 - \alpha \Delta \mathbf{P} - \Delta \mathbf{Q})v = i\partial_t v - (n - 1)v - \partial_t n \left(u_0 + \int_0^t v(s) ds \right).$$

Also the equation above is satisfied in H^{-2} and therefore in the same distributional sense we have

$$\partial_t u = v.$$

Moreover from (3.3) we get

$$\partial_t u = (1 - \alpha \Delta \mathbf{P} - \Delta \mathbf{Q})^{-1} \left(i\partial_t v - (n - 1)v - \partial_t n \left(u_0 + \int_0^t v(s) ds \right) \right) \in \mathcal{C}(I; L^2)$$

therefore $u \in \mathcal{C}^1(I; L^2)$. It is straightforward that $u^\alpha(0, x) = u_0$ and so the proof is complete. \square

Proof of Theorem 3.2. As discussed above, we are going to prove the result by means of a fixed point argument. Let's define the function

$$\tilde{v}(t) := v(t) - U(\alpha t)\mathbf{P}(i\alpha\Delta u_0).$$

We look at the integral formulation for (3.1), namely

$$(3.4) \quad \begin{aligned} v(t) &= U_Z(t)v(0) - i \int_0^t U_Z(t-s)(nv + \partial_t nu)(s) ds \\ n(t) &= \cos(t|\nabla|)n_0 + \frac{\sin(t|\nabla|)}{|\nabla|}n_1 + \int_0^t \frac{\sin((t-s)|\nabla|)}{|\nabla|}\Delta|u|^2 ds, \end{aligned}$$

with u determined by the following elliptic equation

$$-\alpha\nabla \wedge \nabla \wedge u + \nabla \operatorname{div} u = n \left(u_0 + \int_0^t v(s) ds \right) - iv,$$

and $v(0)$ is given by (3.2). This implies that \tilde{v} must satisfy the following integral equation

$$\begin{aligned} \tilde{v}(t) &= U(\alpha t)\mathbf{P}(-in_0u_0) + U(t)\mathbf{Q}(i\Delta u_0 - in_0u_0) \\ &\quad - i \int_0^t U_Z(t-s)(\tilde{v}n + nU(\alpha\cdot)\mathbf{P}(i\alpha\Delta u_0) + \partial_t nu)(s) ds. \end{aligned}$$

Let us consider the space

$$\begin{aligned} X &= \{(\tilde{v}, n) : \tilde{v} \in S^2([0, T]), n \in \mathcal{W}^1([0, T]), \\ &\quad \|\tilde{v}\|_{S^2(I)} \leq M, \|n\|_{\mathcal{W}^1(I)} \leq M\}, \end{aligned}$$

endowed with the norm

$$\|(\tilde{v}, n)\|_X := \|\tilde{v}\|_{S^2(I)} + \|n\|_{\mathcal{W}^1(I)}.$$

Here $0 < T \leq 1$, $M > 0$ will be chosen subsequently and $I := [0, T]$. From the third equation in (3.1) and the definition of \tilde{v} we have

$$(3.5) \quad \begin{aligned} -\alpha\nabla \wedge \nabla \wedge u + \nabla \operatorname{div} u &= -i\tilde{v} - iU(\alpha t)(i\alpha\Delta \mathbf{P}u_0) \\ &\quad - in \left(u_0 + \int_0^t \tilde{v}(s) + U(\alpha s)(i\alpha\Delta \mathbf{P}u_0) ds \right), \end{aligned}$$

thus it is straightforward to see that given n, \tilde{v} , then u is uniquely determined. Furthermore, by applying the projection operators \mathbf{P}, \mathbf{Q} , respectively, to (3.5) we obtain

$$\alpha\Delta \mathbf{P}u = -i\mathbf{P}[\tilde{v} + U(\alpha t)\mathbf{P}(i\alpha\Delta u_0)] + \mathbf{P} \left[n \left(u_0 + \int_0^t \tilde{v}(s) + U(\alpha s)\mathbf{P}(i\alpha\Delta u_0) ds \right) \right]$$

and

$$\Delta \mathbf{Q}u = -i\mathbf{Q}\tilde{v} + \mathbf{Q} \left[n \left(u_0 + \int_0^t \tilde{v}(s) + U(\alpha s)\mathbf{P}(i\alpha\Delta u_0) ds \right) \right].$$

We now estimate the irrotational and solenoidal parts of Δu separately. Let us start with $\mathbf{Q}\Delta u$: by Hölder inequality and Sobolev embedding we obtain

$$\begin{aligned} \|\Delta \mathbf{Q}u\|_{L_t^\infty L_x^2} &\lesssim \|\tilde{v}\|_{L_t^\infty L_x^2} + \|n\|_{L_t^\infty H_x^1} \|u_0\|_{H^2} + T^{1/2} \|n\|_{L_t^\infty H_x^1} \|\tilde{v}\|_{L_t^2 L_x^6} \\ &\quad + T^{1/2} \|n\|_{L_t^\infty H_x^1} \|U(\alpha t)\mathbf{P}(i\alpha\Delta u_0)\|_{L_t^2 L_x^6}. \end{aligned}$$

To estimate the last term, we use the Strichartz estimate in (2.4); let us notice that by choosing the admissible exponents $(q, r) = (2, 6)$ we obtain a factor α^{-1} in the estimate, which balances the term α appearing above. We thus have

$$\|\Delta \mathbf{Q}u\|_{L_t^\infty L_x^2} \lesssim (\|u_0\|_{H^2} + 1)M + M^2.$$

By similar calculations, we also obtain an estimate for $\mathbf{P}\Delta u$,

$$\|\mathbf{P}\Delta u\|_{L_t^\infty L_x^2} \lesssim \|u_0\|_{H^2}^2 + \|u_0\|_{H^2}M + M^2.$$

We then sum up the contributions given by the irrotational and solenoidal parts to get

$$(3.6) \quad \|u\|_{L_t^\infty H_x^2} \lesssim \|u_0\|_{H^2}^2 + \|u_0\|_{H^2}M + M^2 \leq C(\|u_0\|_{H^2})(1 + M^2).$$

Similar calculations also give

$$\begin{aligned} \|u - u'\|_{L^\infty(I; H^2)} &\lesssim \|\tilde{v} - \tilde{v}'\|_{L_t^\infty L_x^2} + \|n - n'\|_{L_t^\infty H_x^1} \\ &\quad + M(\|n - n'\|_{L_t^\infty H_x^1} + \|\tilde{v} - \tilde{v}'\|_{L_t^2 L_x^6}) \\ &\leq C(1 + M)\|(\tilde{v}, n) - (\tilde{v}', n')\|_X. \end{aligned}$$

Given $(\tilde{v}, n) \in X$ we define the map $\Phi : X \rightarrow X$, $\Phi(\tilde{v}, n) = (\Phi_S, \Phi_W)(\tilde{v}, n)$ by

$$\begin{aligned} (3.7) \quad \Phi_S &= U(\alpha t)\mathbf{P}(-in_0u_0) + U(t)\mathbf{Q}(i\Delta u_0 - in_0u_0) \\ &\quad - i \int_0^t U(\alpha(t-s))\mathbf{P}(\tilde{v}n + nU(\alpha\cdot)\mathbf{P}(i\alpha\Delta u_0) + \partial_t nu)(s) ds \\ &\quad - i \int_0^t U(t-s)\mathbf{Q}(n\tilde{v} + nU(\alpha\cdot)(i\alpha\Delta u_0) + \partial_t nu)(s) ds \\ (3.8) \quad \Phi_W &= \cos(t|\nabla|)n_0 + \frac{\sin(t|\nabla|)}{|\nabla|}n_1 + \int_0^t \frac{\sin((t-s)|\nabla|)}{|\nabla|}\Delta|u|^2(s) ds, \end{aligned}$$

where u in the formulas above is given by (3.5) and its $L_t^\infty H_x^2$ norm is bounded in (3.6). Let us first prove that, by choosing T and M properly, Φ maps X into itself.

Let us first analyze the Schrödinger part (3.7), by the Strichartz estimates in Lemma 2.2, Hölder inequality and Sobolev embedding we have

$$\|U(\alpha t)\mathbf{P}(-in_0u_0) + U(t)\mathbf{Q}(i\Delta u_0 - in_0u_0)\|_{L^q L^r} \lesssim \|u_0\|_{H^2} + \|n_0\|_{H^1}\|u_0\|_{H^2}$$

We treat the inhomogenous part similarly,

$$\begin{aligned} &\left\| \int_0^t U_Z(t-s)(n\tilde{v} + nU(\alpha s)(i\alpha\mathbf{P}\Delta u_0))(s) ds \right\|_{L_t^q L_x^r} \lesssim \|n\tilde{v} + nU(\alpha\cdot)(i\alpha\mathbf{P}\Delta u_0)\|_{L_t^1 L_x^2} \\ &\lesssim T^{1/2}\|n\|_{L_t^\infty H_x^1}(\|\tilde{v}\|_{L_t^2 L_x^6} + \|U(\alpha t)\mathbf{P}(i\alpha\Delta u_0)\|_{L_t^2 L_x^6}) \lesssim T^{1/2}M(M + \|u_0\|_{H^2}). \end{aligned}$$

where in the last inequality we again used (2.4) with $(q, r) = (2, 6)$. Similarly,

$$\begin{aligned} &\left\| \int_0^t U_Z(t-s)(\partial_t nu)(s) ds \right\|_{L_t^q L_x^r} \lesssim T\|\partial_t n\|_{L_t^\infty L_x^2}\|u\|_{L_t^\infty H_x^2} \\ &\lesssim C(\|u_0\|_{H^2})TM(1 + M^2), \end{aligned}$$

where in the last line we use the bound (3.6). Collecting these estimates we get

$$(3.9) \quad \|\Phi_S(\tilde{v}, n)\|_{L_t^q L_x^r} \leq C(\|u_0\|_{H^2}, \|n_0\|_{L^2}) + CT^{1/2}M(1 + M).$$

For the wave component we use formula (3.8) and Hölder inequality to obtain

$$\begin{aligned}\|\Phi_W(v, n)\|_{\mathcal{W}^1(I)} &\leq C(1+T)\|n_0\|_{H^1} + \|n_1\|_{L^2} + \|\Delta|u|^2\|_{L_t^1 L_x^2} \\ &\leq C(\|n_0\|_{H^1} + \|n_1\|_{L^2}) + T\|u\|_{L_t^\infty H_x^2}^2,\end{aligned}$$

where we used the fact that $H^2(\mathbb{R}^3)$ is an algebra. From (3.6) we infer

$$(3.10) \quad \|\Phi_W(v, n)\|_{\mathcal{W}^1(I)} \leq C(\|n_0\|_{H^1}, \|n_1\|_{L^2}) + T(M + M^4).$$

The bounds (3.9) and (3.10) together yield

$$\|\Phi(\tilde{v}, n)\|_X \leq C(\|(u_0, n_0, n_1)\|_{\mathcal{H}_2}) + CT^{1/2}M(1 + M^3).$$

Let us choose M such that

$$\frac{M}{2} = C(\|(u_0, n_0, n_1)\|_{\mathcal{H}_2})$$

and T such that

$$CT^{1/2}(1 + M^3) < \frac{1}{2},$$

we then obtain $\|\Phi(\tilde{v}, n)\|_X \leq M$. Hence Φ maps X into itself. It thus remains to prove that Φ is a contraction. Arguing similarly to what we did before we obtain

$$\begin{aligned}\|\Phi_S(\tilde{v}, n) - \Phi_S(\tilde{v}', n')\|_{L_t^q L_x^r} &\leq CT^{1/2}(1 + M)\|(\tilde{v}, n) - (\tilde{v}', n')\|_{L_t^q L_x^r} \\ \|\Phi_W(\tilde{v}, n) - \Phi_W(\tilde{v}', n')\|_{\mathcal{W}^1(I)} &\leq CT(1 + M^3)\|(\tilde{v}, n) - (\tilde{v}', n')\|_{\mathcal{W}^1(I)}.\end{aligned}$$

By possibly choosing a smaller $T > 0$ such that $CT^{1/2}(1 + M^3) < 1$ then we see that $\Phi : X \rightarrow X$ is a contraction and consequently there exists a unique $(\tilde{v}, n) \in X$ which is a fixed point for X . Let us notice that the time T depends only on M , hence $T = T(\|(u_0, n_0, n_1)\|_{\mathcal{H}_2})$. Furthermore from the definition of \tilde{v} it follows that (u, v, n) is a solution to (3.1), where $v = \tilde{v} + U(\alpha t)\mathbf{P}(i\alpha\Delta u_0)$. From (3.6) we also see that the $L_t^\infty H_x^2$ norm of u is uniformly bounded in α .

Finally, from standard arguments we extend the solution on a maximal time interval, on which the standard blow-up alternative holds true and we can also infer the continuous dependence on the initial data. \square

4. CONVERGENCE OF SOLUTIONS.

In this Section we study the electrostatic limit for the vectorial Zakharov system (1.1). As stated in Theorem 1.3 we will show that solutions (u^α, n^α) to (1.1) converge in the space \mathcal{X}_T towards solutions (u^∞, n^∞) to (1.2), for any $0 < T < T_{max}^\infty$.

As in previous Section we are going to study the convergence through the two associated auxiliary systems, namely (3.1) for (1.1) and

$$(4.1) \quad \begin{cases} i\partial_t v^\infty + \Delta v^\infty = \mathbf{Q}(n^\infty v^\infty + \partial_t n^\infty u^\infty) \\ \partial_{tt} n^\infty - \Delta n^\infty = \Delta|u^\infty|^2 \\ iv^\infty + \Delta u^\infty = \mathbf{Q}(n^\infty u^\infty) \end{cases}$$

for (1.2). Again, we use those auxiliary systems in order to overcome the difficulty generated by the loss of derivatives on the terms $|u^\alpha|^2$ and $|u^\infty|^2$. On the other hand this introduces a new difficulty as the initial datum for v^α ,

$$(4.2) \quad v^\alpha(0) = -i\alpha\nabla \wedge \nabla \wedge u_0 + i\nabla \operatorname{div} u_0 - in_0 u_0,$$

is not uniformly bounded in L^2 for every $\alpha \geq 1$. Even more, the term

$$\mathbf{P}U_Z(t)v(0) = U(\alpha t)\mathbf{P}(i\alpha\Delta u_0 - in_0 u_0)$$

does not converge to zero, in general. Indeed, by (2.4) the best estimate in α for this term would be by estimating it in $L_t^2 L_x^6$, for which we have

$$\|U(\alpha t)\mathbf{P}(i\alpha\Delta u_0 - in_0 u_0)\|_{L_t^2 L_x^6} \lesssim \|\mathbf{P}\Delta u_0\|_{L^2} + \alpha^{-1}\|n_0\|_{H^1}\|u_0\|_{H^2}.$$

For this reason we will need a set of well-prepared initial data for (3.1). More specifically we consider $(u_0^\alpha, n_0^\alpha, n_1^\alpha) \in \mathcal{H}_2$ such that

$$(4.3) \quad \|(u_0^\alpha, n_0^\alpha, n_1^\alpha) - (u_0^\infty, n_0^\infty, n_1^\infty)\|_{\mathcal{H}_2} \rightarrow 0$$

for some $(u_0^\infty, n_0^\infty, n_1^\infty) \in \mathcal{H}_2$ and

$$(4.4) \quad \|\mathbf{P}u_0^\alpha\|_{H^2} \rightarrow 0.$$

This clearly implies that the initial datum for the limit equation (1.2) is irrotational, i.e. $\mathbf{P}u_0^\infty = 0$.

To prove the convergence result stated in Theorem 1.3 we will study the convergence from (3.1) to (4.1). The main result of this Section is the following.

Theorem 4.1. *Let $\alpha \geq 1$ and let $(u_0^\alpha, n_0^\alpha, n_1^\alpha), (u_0^\infty, n_0^\infty, n_1^\infty) \in \mathcal{H}_2$ be initial data such that (4.3) and (4.4) hold true. Let $(u^\alpha, v^\alpha, n^\alpha)$ be the maximal solution to (3.1) with Cauchy data $(u_0^\alpha, n_0^\alpha, n_1^\alpha)$ given by Theorem 3.2 and analogously let $(u^\infty, v^\infty, n^\infty)$ be the maximal solution to (4.1) in the interval $[0, T_{max}^\infty)$ accordingly to Theorem 1.1. Then for any $0 < T < T_{max}^\infty$ we have*

$$\lim_{\alpha \rightarrow \infty} \|(u^\alpha, v^\alpha, n^\alpha) - (u^\infty, v^\infty, n^\infty)\|_{L^\infty(0, T; \mathcal{H}_2)} = 0.$$

The proof of the Theorem above is divided in two main steps. First of all we prove in Lemma 4.2 that, as long as the \mathcal{H}_2 norm of $(u^\alpha(T), n^\alpha(T), \partial_t n^\alpha(T))$ is bounded, then the convergence holds true in $[0, T]$. The second one consists in proving that the \mathcal{H}_2 bound on $(u^\alpha(T), n^\alpha(T), \partial_t n^\alpha(T))$ holds true for any $0 < T < T_{max}^\infty$. A similar strategy of proof is already exploited in the literature to study the asymptotic behavior of time oscillating nonlinearities, see for example [CS] where the authors consider a time oscillating nonlinearity or [AW] where in a system of two nonlinear Schrödinger equations a rapidly varying linear coupling term is averaging out the effect of nonlinearities. We also mention [CPS] where a similar strategy is also used to study a time oscillating critical Korteweg-de Vries equation.

Lemma 4.2. *Let $(u^\alpha, v^\alpha, n^\alpha), (u^\infty, v^\infty, n^\infty)$ be defined as in the statement of Theorem 4.1 and let us assume that for some $0 < T_1 < T_{max}^\infty$ we have*

$$\sup_{\alpha \geq 1} \|(u^\alpha, n^\alpha, \partial_t n^\alpha)\|_{L^\infty(0, T_1; \mathcal{H}_2)} < \infty.$$

It follows that

$$\lim_{\alpha \rightarrow \infty} \left(\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} \right) = 0,$$

where all the norms are taken in the space-time slab $[0, T_1] \times \mathbb{R}^3$. In particular we have

$$\lim_{\alpha \rightarrow \infty} \|(u^\alpha, n^\alpha, \partial_t n^\alpha) - (u^\infty, n^\infty, \partial_t n^\infty)\|_{L^\infty(0, T_1; \mathcal{H}_2)} = 0.$$

We assume for the moment that Lemma 4.2 holds true, then we first show how this implies Theorem 4.1.

Proof of Theorem 4.1. Let $0 < T < T_{max}^\infty$ be fixed and let us define

$$N := 2\|(u^\infty, n^\infty, \partial_t n^\infty)\|_{L^\infty(0, T; \mathcal{H}_2)}.$$

From the local well-posedness theory, see Proposition 3.1, there exists $\tau = \tau(N)$ such that the solution $(u^\alpha, n^\alpha, \partial_t n^\alpha)$ to (3.1) exists on $[0, \tau]$ and we have

$$\|(u^\alpha, n^\alpha, \partial_t n^\alpha)\|_{L^\infty(0, T_1; \mathcal{H}_2)} < \infty.$$

We observe that, because of what we said before, the choice $T_1 = \tau$ is always possible. By the Lemma 4.2 we infer that

$$\lim_{\alpha \rightarrow \infty} \|(u^\alpha, n^\alpha, \partial_t n^\alpha) - (u^\infty, n^\infty, \partial_t n^\infty)\|_{L^\infty(0, T_1; \mathcal{H}_2)} = 0.$$

On the other hand by the definition of N we have that, for $\alpha \geq 1$ large enough,

$$\begin{aligned} \|(u^\alpha, n^\alpha, \partial_t n^\alpha)(T_1)\|_{\mathcal{H}_2} &\leq \|(u^\alpha, n^\alpha, \partial_t n^\alpha)(T_1) - (u^\infty, n^\infty, \partial_t n^\infty)(T_1)\|_{\mathcal{H}_2} \\ &\quad + \|(u^\infty, n^\infty, \partial_t n^\infty)(T_1)\|_{\mathcal{H}_2} \leq N. \end{aligned}$$

Consequently we can apply Proposition 3.1 to infer that (u^α, n^α) exists on a larger time interval $[0, T_1 + \tau]$, provided $T_1 + \tau \leq T$, and again

$$\|(u^\alpha, n^\alpha, \partial_t n^\alpha)\|_{L^\infty(0, T_1 + \tau; \mathcal{H}_2)} \leq 2N.$$

We can repeat the argument iteratively on the whole interval $[0, T]$ to infer

$$\|(u^\alpha, n^\alpha, \partial_t n^\alpha)\|_{L^\infty(0, T; \mathcal{H}_2)} \leq 2N.$$

By using Lemma 4.2 this proves the Theorem. \square

It only remains now to prove Lemma 4.2.

Proof of Lemma 4.2. Let us fix

$$M := \sup_{\alpha} \sup_{[0, T_1]} \|(u^\alpha, n^\alpha, \partial_t n^\alpha)(t)\|_{\mathcal{H}_2}.$$

By using the integral formulation for (3.1) and (4.1) we have

$$\begin{aligned} v^\alpha(t) - v^\infty(t) &= U(\alpha t) \mathbf{P}(\alpha \Delta u_0^\alpha - i u_0^\alpha n_0^\alpha) + U(t) \mathbf{Q}(v_0^\alpha - v_0^\infty) \\ &\quad - i \int_0^t U(\alpha(t-s)) [\mathbf{P}(\partial_t(n^\alpha u^\alpha))](s) ds \\ &\quad - i \int_0^t U(t-s) [\mathbf{Q}(\partial_t(n^\alpha u^\alpha) - \partial_t(n^\infty u^\infty))](s) ds. \end{aligned}$$

Now we use the Strichartz estimates in Lemma 2.2 to get

$$\begin{aligned} \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} &\lesssim \|\mathbf{P} u_0^\alpha\|_{H^2} + \alpha^{-1} \|n_0^\alpha\|_{H^1} \|u_0^\alpha\|_{H^2} + \|v_0^\alpha - v_0^\infty\|_{L^2} \\ &\quad + \alpha^{-1/2} \|n^\alpha v^\alpha + \partial_t n^\alpha u^\alpha\|_{L_t^1 L_x^2} \\ &\quad + \|n^\alpha v^\alpha - n^\infty v^\infty\|_{L_t^1 L_x^2} + \|\partial_t n^\alpha u^\alpha - \partial_t n^\infty u^\infty\|_{L_t^1 L_x^2}. \end{aligned}$$

It is straightforward to check that, by Hölder inequality and Sobolev embedding,

$$\begin{aligned} \|n^\alpha v^\alpha + \partial_t n^\alpha u^\alpha\|_{L_t^1 L_x^2} &\leq C(T, M), \\ \|n^\alpha v^\alpha - n^\infty v^\infty\|_{L_t^1 L_x^2} &\lesssim T^{1/2} (\|n^\alpha - n^\infty\|_{L_t^\infty H_x^1} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6}), \\ \|\partial_t n^\alpha u^\alpha - \partial_t n^\infty u^\infty\|_{L_t^1 L_x^2} &\lesssim T (\|\partial_t n^\alpha - \partial_t n^\infty\|_{L_t^\infty L_x^2} + \|u^\alpha - u^\infty\|_{L_t^\infty H_x^2}). \end{aligned}$$

By putting all the estimates together we obtain

$$\begin{aligned} \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} &\lesssim \|\mathbf{P}u_0^\alpha\|_{H^2} + \alpha^{-1}\|n_0^\alpha\|_{H^1}\|u_0^\alpha\|_{H^2}\|u_0^\alpha - u_0^\infty\|_{H^2} + \alpha^{-1/2} + \|n_0^\alpha - n_0^\infty\|_{H^1} \\ &\quad + T^{1/2}(\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1}) \end{aligned}$$

To estimate the wave part in (3.1) and (4.1), we write

$$\begin{aligned} n^\alpha - n^\infty &= \cos(t|\nabla|)(n_0^\alpha - n_0^\infty) - \frac{\sin(t|\nabla|)}{|\nabla|}(n_1^\alpha - n_1^\infty) \\ &\quad + \int_0^t \frac{\sin((t-s)|\nabla|)}{|\nabla|} \Delta(|u^\alpha|^2 - |u^\infty|^2) ds, \end{aligned}$$

whence, by using again that $H^2(\mathbb{R}^3)$ is an algebra,

$$\|n^\alpha - n^\infty\|_{\mathcal{W}^1} \lesssim \|n_0^\alpha - n_0^\infty\|_{H^1} + \|n_1^\alpha - n_1^\infty\|_{L^2} + T\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2}.$$

The estimate for the difference $u^\alpha - u^\infty$ is more delicate. From the third equations in (3.1) and (4.1) we have

$$-\alpha \nabla \wedge \nabla \wedge u^\alpha + \nabla \operatorname{div}(u^\alpha - u^\infty) = i(v^\alpha - v^\infty) - n^\alpha u^\alpha + \mathbf{Q}(n^\infty u^\infty).$$

Again, here we estimate separately the irrotational and solenoidal parts of the difference. For the solenoidal part we obtain

$$\alpha \|\mathbf{P} \Delta u^\alpha\|_{L_t^\infty L_x^2} \lesssim \|v^\alpha\|_{L_t^\infty L_x^2} + \|n^\alpha u^\alpha\|_{L_t^\infty L_x^2}.$$

To estimate the $L_t^\infty L_x^2$ norm of v^α on the right hand side we use (3.4) and Strichartz estimates to infer

$$\|v^\alpha\|_{L_t^\infty L_x^2} \lesssim \alpha \|\mathbf{P}u_0^\alpha\|_{H^2} \|u_0^\alpha\|_{H^2} \|n_0^\alpha\|_{H^1} + 1.$$

Hence

$$\alpha \|\mathbf{P} \Delta u^\alpha\|_{L_t^\infty L_x^2} \lesssim \alpha \|\mathbf{P}u_0^\alpha\|_{H^2} + \|u_0^\alpha\|_{H^2} \|n_0^\alpha\|_{H^1} + 1.$$

For the irrotational part

$$(4.5) \quad \|\mathbf{Q} \Delta(u^\alpha - u^\infty)\|_{L_t^\infty L_x^2} \lesssim \|\mathbf{Q}(v^\alpha - v^\infty)\|_{L_t^\infty L_x^2} + \|n^\alpha - u^\alpha - n^\infty u^\infty\|_{L_t^\infty L_x^2}.$$

By using (3.4), the analogue integral formulation for v^∞ and by applying the Helmholtz projection operator \mathbf{Q} to their difference we have that the first term on the right hand side is bounded by

$$\begin{aligned} \|\mathbf{Q}(v^\alpha - v^\infty)\|_{L_t^\infty L_x^2} &\lesssim \|u_0^\alpha - u_0^\infty\|_{H^2} + \|n_0^\alpha - n_0^\infty\|_{H^1} \\ &\quad + T^{1/2} \left(\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} \right) \end{aligned}$$

The second term on the right hand side of (4.5) is estimated by

$$\begin{aligned} \|n^\alpha u^\alpha - n^\infty u^\infty\|_{L_t^\infty L_x^2} &\lesssim \|n^\alpha - n^\infty\|_{L_t^\infty L_x^2} \|u^\alpha\|_{L_t^\infty H_x^2} \\ &\quad + \|n^\infty(u_0^\alpha - u_0^\infty)\|_{L_t^\infty L_x^2} + \left\| n^\infty \int_0^t (v^\alpha - v^\infty) \right\|_{L_t^\infty L_x^2} \\ &\lesssim (\|n_0^\alpha - n_0^\infty\|_{L^2} + T\|\partial_t n^\alpha - \partial_t n^\infty\|_{L_t^\infty L_x^2}) M \\ &\quad + \|n^\infty\|_{L_t^\infty L_x^2} \|u_0^\alpha - u_0^\infty\|_{H_x^2} \\ &\quad + T^{1/2} \|n^\infty\|_{L_t^\infty H_x^1} \|v^\alpha - v^\infty\|_{L_t^2 L_x^6}. \end{aligned}$$

By summing up the two contribution in (4.5) we then get

$$\begin{aligned} \|\mathbf{Q}\Delta(u^\alpha - u^\infty)\|_{L_t^\infty L_x^2} &\lesssim \|u_0^\alpha - u_0^\infty\|_{H^2} + \|n_0^\alpha - n_0^\infty\|_{H^1} \\ &\quad + T^{1/2} \left(\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} \right). \end{aligned}$$

Finally, we notice that, by using the Schrödinger equations in (1.1) and (1.2), we have

$$\|u^\alpha - u^\infty\|_{L_t^\infty L_x^2} \lesssim T \left(\|n^\alpha - n^\infty\|_{L_t^\infty H_x^1} + \|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} \right),$$

so that

$$\begin{aligned} \|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} &\lesssim \|u_0^\alpha - u_0^\infty\|_{H^2} + \|n_0^\alpha - n_0^\infty\|_{H^1} + \|\mathbf{P}u_0^\alpha\|_{H^2} + \alpha^{-1} \\ &\quad + T^{1/2} \left(\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} \right) \end{aligned}$$

Now we put everything together, we finally obtain

$$\begin{aligned} \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} + \|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} &\lesssim \\ &\lesssim \|\mathbf{P}u_0^\alpha\|_{H^2} + \alpha^{-1} + \|u_0^\alpha - u_0^\infty\|_{H^2} + \|n_0^\alpha - n_0^\infty\|_{H^1} + \|n_1^\alpha - n_1^\infty\|_{L^2} \\ &\quad + T^{1/2} \left(\|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} + \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} \right) \end{aligned}$$

By choosing T small enough depending on M we can infer

$$\begin{aligned} \|v^\alpha - v^\infty\|_{L_t^2 L_x^6} + \|n^\alpha - n^\infty\|_{\mathcal{W}^1} + \|u^\alpha - u^\infty\|_{L_t^\infty H_x^2} &\lesssim \\ &\lesssim \|\mathbf{P}u_0^\alpha\|_{H^2} + \alpha^{-1} + \|u_0^\alpha - u_0^\infty\|_{H^2} + \|n_0^\alpha - n_0^\infty\|_{H^1} + \|n_1^\alpha - n_1^\infty\|_{L^2}. \end{aligned}$$

This proves the convergence in the time interval $[0, T]$, for $T > 0$ small enough. Let now $0 < T_1$ be as in the statement of Lemma, we can divide $[0, T_1]$ into many subintervals of length T such that the convergence holds in any small interval. By gluing them together we prove the Lemma. \square

REFERENCES

- [AA] ADDED H., ADDED S., *Equation of Langmuir turbulence and nonlinear Schrödinger equation: smoothness and approximation*, J. Funct. Anal. 79 (1988), 183–210.
- [ASS] ANTONELLI P., SAUT J.-C., SPARBER C., *Well-Posedness and averaging of NLS with time-periodic dispersion management*, Adv. Diff. Eqns. 18, no. 1/2 (2013), 49–68.
- [AW] ANTONELLI P., WEISHAEUPL R., *Asymptotic Behavior of Nonlinear Schrödinger Systems with Linear Coupling*, JHDE 11, no. 1 (2014), 159–183.
- [BC] BOURGAIN, J., COLLIANDER, J., *On wellposedness of the Zakharov system*, International Mathematics Research Notices, 1996, no. 11;
- [BH] BEJENARU, I. HERR, S., *Convolutions of singular measures and applications to the Zakharov system*, J. Funct. Anal. 261 (2011), no. 2, 478–506;
- [CPS] CARVAJAL X., PANTHEE M., SCIALOM M., *On the critical KdV equation with time-oscillating nonlinearity*, Diff. Int. Eqns. 24, no. 5/6 (2011), 541–567.
- [C] CAZENAVE, T., *Semilinear Schrödinger equations*, Courant Lecture Notes in Mathematics, 10. Ams;
- [CS] CAZENAVE T., SCIALOM M. *A Schrödinger equation with time-oscillating nonlinearity*, Rev. Mat. Comp. 23, no. 2 (2010), 321–339.
- [CW] CAZENAVE, T., WEISSLER, F.B., *The Cauchy problem for the nonlinear Schrödinger equation in H^1* , Manuscripta Math. 61 (1988), no. 4, 477–494;
- [CCS] COLLIANDER, J., CZUBAK, M., SULEM, C., *Lower bound for the rate of blow-up of singular solutions of the Zakharov system in \mathbb{R}^3* , Jour. Hyperbolic Differ. Equ. 10 (2013), no. 3, 523–536;

- [G] GALUSINSKI C., *A singular perturbation problem in a system of nonlinear Schrödinger equations occurring in a Langmuir turbulence*, Rapport interne 98011 (1998), Mathématiques Appliquées de Bordeaux.
- [GM] GLANGETAS, L., MERLE, F., *Existence of self-similar blowup solutions for Zakharov equations in dimension 2. Part I and II*, Commun. Math. Phys. 160, 173-215 (1994);
- [GTV] GINIBRE, J., TSUTSUMI, Y., VELO, G., *On the Cauchy problem for the Zakharov system* J. Funct. Anal. 151 (1997), 384-436;
- [GV] GINIBRE, J., VELO, G., *The global Cauchy problem for the nonlinear Schrödinger equation revisited*, Ann. Inst. H. Poincaré, Anal. Non Linéaire 2: 309-327 (1985);
- [KPV] KENIG, C. E., PONCE, G. AND VEGA, L., *On the Zakharov and Zakharov-Schulman systems*, J. Funct. Anal. 127 (1995) no. 1, 204-234;
- [KT] KEEL, M., TAO, T., *Endpoint Strichartz inequalities*, Amer. Jour. Math. 120 (1988) 955-980;
- [LP] LINARES, F., PONCE, G., *Introduction to nonlinear dispersive equations*, Universitext. Springer, New York, 2009;
- [M] MERLE, F., *Blow-up results of Viriel Type for Zakharov System*, Comm. in Math. Ph. 175 (1996), 433-455;
- [MN] MASMOUDI, N., NAKANISHI, K., *Energy convergence for singular limits of Zakharov type systems*, Invent. Math. 172 (2008), no. 3, 535-583;
- [OT1] OZAWA, T., TSUTSUMI, Y., *Existence of smoothing effect of solutions for the Zakharov equations*, Publ. Res. Inst. Math. Sci. 28 (1992), no. 3, 329-361;
- [OT2] OZAWA, T., TSUTSUMI, Y., *Global existence and asymptotic behavior of solutions for the Zakharov equations in three space dimensions*, Adv. Math. Sci. Appl. 3 (1993/94), Special Issue, 301-334;
- [SW] SCHOCHET, S.H., WEINSTEIN, M.I., *The nonlinear Schrödinger limit of the Zakharov equations governing Langmuir turbulence*, Comm. Math. Phys., 106 (1986), 569-580;
- [STR] STRICHARTZ, R., *Restrictions of Fourier transforms to quadratic surfaces and decay of solutions of wave equations*, Duke Math. Jour. 44 (1977), no. 3, 705-714;
- [SS] SULEM, C., SULEM, P.L., *The Nonlinear Schrödinger Equation: Self-Focusing and Wave Collapse*, Applied Mathematical Sciences, 139. Springer-Verlag, New York, 1999;
- [SS2] C. Sulem, P.L. Sulem, *Quelques résultats de régularité pour les équations de la turbulence de Langmuir*, C. R. Acad. Sci. Paris **289** (1979), 173-176.
- [T] TEXIER, B., *Derivation of the Zakharov equations*, Arch. Ration. Mech. Anal. 184 (2007), no. 1, 121-183;
- [TtH] THORNHILL S.G., TER HAAR D., *Langmuir Turbulence and Modulational Instability*, Phys. Reports 43 (1978), 43-99;
- [Y] YAJIMA, K., *Existence of solutions for Schrödinger evolution equations*, Comm. Math. Phys. 110 (1987), no. 3, 415-426;
- [Z] ZAKHAROV, V.E., *Collapse of Langmuir waves*, Sov. Phys. JETP, 35 (1972), 908-914.

GSSI, GRAN SASSO SCIENCE INSTITUTE, VIALE F. CRISPI 7, 67100 L'AQUILA, ITALY
E-mail address: `paolo.antonelli@gssi.infn.it`

SCUOLA NORMALE SUPERIORE, PIAZZA DEI CAVALIERI, 7, 56126 PISA, ITALY
E-mail address: `luigi.forcella@sns.it`